

# Time-Resolved Fourier Transform Infrared Spectroscopy of the Polarizable Proton Continua and the Proton Pump Mechanism of Bacteriorhodopsin

Jianping Wang and Mostafa A. El-Sayed

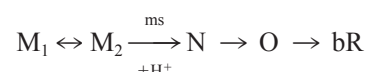
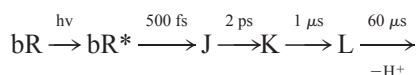
Laser Dynamics Laboratory, School of Chemistry and Biochemistry, Georgia Institute of Technology, Atlanta, Georgia 30332 USA

**ABSTRACT** Nanosecond-to-microsecond time-resolved Fourier transform infrared (FTIR) spectroscopy in the 3000–1000-cm<sup>-1</sup> region has been used to examine the polarizable proton continua observed in bacteriorhodopsin (bR) during its photocycle. The difference in the transient FTIR spectra in the time domain between 20 ns and 1 ms shows a broad absorption continuum band in the 2100–1800-cm<sup>-1</sup> region, a bleach continuum band in the 2500–2150-cm<sup>-1</sup> region, and a bleach continuum band above 2700 cm<sup>-1</sup>. According to Zundel (G. Zundel, 1994, *J. Mol. Struct.* 322:33–42), these continua appear in systems capable of forming polarizable hydrogen bonds. The formation of a bleach continuum suggests the presence of a polarizable proton in the ground state that changes during the photocycle. The appearance of a transient absorption continuum suggests a change in the polarizable proton or the appearance of new ones. It is found that each continuum has a rise time of less than 80 ns and a decay time component of ~300 μs. In addition, it is found that the absorption continuum in the 2100–1800-cm<sup>-1</sup> region has a slow rise component of 190 ns and a fast decay component of ~60 μs. Using these results and those of the recent x-ray structural studies of bR<sub>570</sub> and M<sub>412</sub> (H. Luecke, B. Schobert, H.T. Richter, J.-P. Cartailier, and J. K. Lanyi, 1999, *Science* 286:255–260), together with the already known spectroscopic properties of the different intermediates in the photocycle, the possible origins of the polarizable protons giving rise to these continua during the bR photocycle are proposed. Models of the proton pump are discussed in terms of the changes in these polarizable protons and the hydrogen-bonded chains and in terms of previously known results such as the simultaneous deprotonation of the protonated Schiff base (PSB) and Tyr185 and the disappearance of water molecules in the proton release channel during the proton pump process.

## INTRODUCTION

Bacteriorhodopsin (bR), present in the purple membrane of *Halobacterium salinarum*, functions as a proton pump upon visible light absorption (Stoeckenius and Lozier, 1974; Birge, 1981; Mathies et al., 1991; Rothschild, 1992; El-Sayed, 1992; Ebrey, 1993; Lanyi, 1993). There are 248 amino acid residues in the single polypeptide chain, arranged in seven α-helices (Henderson, 1979). The three-dimensional structure of bR has been determined at high resolution by a number of researchers, using cryoelectron microscopy (Henderson et al., 1990) and x-ray crystallography (Pebay-Peyroula et al., 1997; Luecke et al., 1998; Essen et al., 1998). The retinal chromophore that is covalently bound to Lys216 through a protonated Schiff base (PSB) undergoes a photoisomerization from all-*trans* to 13-*cis* and triggers a cyclic thermal reaction that involves a series of photointermediates, J, K, L, M<sub>1</sub>, M<sub>2</sub>, N, and O, with lifetimes ranging from half a picosecond to tens of

milliseconds:



The nature of J is under discussion, and recent results suggest it is an excited-state conformer of all-*trans* retinal (Atkinson et al., 2000). Retinal, protein residue, and hydrogen bonding structural changes are expected to be associated with the bR photocycle. Upon retinal *cis-trans* photoisomerization and K formation, there are changes in the protein peptide backbone vibration frequency (Earnest et al., 1986), in the hydrogen-bonded network near the retinal (Kandori et al., 1998), as well as in the O-H stretching vibration of a water molecule as shown by using trapping technique and difference Fourier transform infrared (FTIR) spectroscopy at near 80 K (Maeda et al., 1992; Fischer et al., 1994). Upon M formation, the difference FTIR spectrum in the amide I and amide II bands suggest large changes in the protein secondary structure involving one or more helices (Rothschild, 1992). Room-temperature, time-resolved FTIR reveals protein structural changes that are partially blocked at low temperature, such as a greater degree of deprotonation of Asp96 in the L intermediate (Chen and Braiman, 1991). Large changes in the secondary and tertiary structures of bR upon N formation have been reported by using

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Address reprint requests to Dr. Mostafa A. El-Sayed, Laser Dynamics Laboratory, School of Chemistry and Biochemistry, Georgia Institute of Technology, Atlanta, GA 30332-0400. Tel.: 404-894-0292; Fax: 404-894-0294; E-mail: Mostafa.El-Sayed@chemistry.gatech.edu.

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neutron and x-ray diffraction techniques (Dencher et al., 1990; Papadopoulos et al., 1990; Koch et al., 1991; Oka et al., 1999) and by using FTIR difference spectroscopy (Rothschild et al., 1990b; Nilsson et al., 1995). In the L-M transition, it is now established that a proton is transferred from the PSB to Asp85 (Braiman et al., 1988) and another proton appears on the extracellular surface (Heberle and Dencher, 1992; Alexiev et al., 1995). The origin of the latter (and thus the mechanism of the proton pump) has been the subject of intense research recently and has been proposed to be or involve several groups, such as Arg82 (Otto et al., 1990; Balashov et al., 1993, 1995; Govindjee et al., 1996), Glu194 (Dioumaev et al., 1998), or Glu204 (Misra et al., 1997; Kandori et al., 1997; Richter et al., 1996; Brown et al., 1995, 1996). Recently, it has been proposed to be due to the polarizable hydrogen-bonded proton with the spectroscopic signature in the region of 1900–1800  $\text{cm}^{-1}$  (Rammelsberg et al., 1998; Wang and El-Sayed, 2000). Earlier, an ionized water molecule within the active site was also proposed to be the source of this proton (El-Sayed et al., 1995). Very recently, the x-ray structural changes upon the formation of M in a D96N mutant has recently been reported (Luecke et al., 1999a), and a proposed proton pump mechanism has been suggested, which contradicts the FTIR results (Rammelsberg et al., 1998; Zscherp et al., 1999).

Thus, despite the great amount of structural information about bR and its deprotonated form, the molecular origin of the proton pumped to the surface and the details of the proton pump mechanism itself remains very much unsolved. It is hoped that with more time-resolved spectral studies together with theoretical calculations and the determination of the x-ray structure of the other intermediates, a final accurate picture of the proton pump process can be determined. The present work is an effort in this direction and is aimed at a detailed study of the time-resolved infrared continua spectra believed to be due to the changes in the polarizable hydrogen-bonded protons in bR and its intermediates.

By using model compounds, it has been suggested (Merz and Zundel, 1986) that the hydrogen bonds between a Schiff base and a carboxylic acid may show large proton polarizability and, therefore, give rise to several very broad absorption continua in the infrared frequency region of 2800–800  $\text{cm}^{-1}$  (Zundel and Merz, 1984; Mueller et al., 1991; Olejnik et al., 1992). The difference infrared continua observed for low-temperature trapped intermediates between bR and M have been reported and discussed previously (Olejnik et al., 1992; Zundel, 1994). The assignment of the transient continuum observed in the region of 1900–1800  $\text{cm}^{-1}$  to the polarizable proton being pumped to the extracellular surface has been suggested recently by time-resolved FTIR studies (Riesle et al., 1996; Rammelsberg et al., 1998), and was more recently supported by the correspondence between the kinetics of the continuum and that of

M formation, as well as the isotope effect on the transient continuum spectrum (Wang and El-Sayed, 2000).

In the present study, we examine the time-resolved infrared continuum spectra in the 3000–1800- $\text{cm}^{-1}$  region for different intermediates formed during the bR photocycle at physiological temperature in the tens of nanoseconds to the hundreds of microseconds time domain. The results suggest that the ground state of native bR has a polarizable hydrogen-bonded network that changes its nature upon the formation of K, leading to an observed time-dependent bleach continua in the 2500–2150- $\text{cm}^{-1}$  region and above 2700  $\text{cm}^{-1}$ . As a result, a new transient absorption continuum appears in the 2100–1800- $\text{cm}^{-1}$  region. This transient absorption continuum band decays during the M formation (in  $\sim 60 \mu\text{s}$ ), whereas the ground-state polarizable hydrogen bond recovers on a longer time scale ( $\sim 300 \mu\text{s}$ ). The observed results, combined with previously published optical and FTIR spectroscopic observations as well as the results of the recent x-ray studies are all used to discuss the nature of these polarizable proton bands and the possible proton pump mechanisms in bR.

## MATERIALS AND METHODS Sample

Purple membrane was isolated from the *Halobacterium salinarum* strain ET1001 according to a general procedure (Stoeckenius and Lozier, 1974). Purple membrane fragments containing bR molecules are suspended in  $\text{H}_2\text{O}$  at pH 7. Concentrated bR samples are squeezed in between two  $\text{CaF}_2$  plates to give an optical density at protein amide I and II bands (1700–1500  $\text{cm}^{-1}$ ) of  $\sim 0.5$ – $1.0$ .

## Time-resolved FTIR spectroscopy

The step-scan FTIR measurements were carried out on the Nicolet Magna-IR 860 spectrometer. A 50-MHz mercury-cadmium-telluride (MCT) detector (Kolmar Technologies, Inc., Congers, GA) was used, with appropriate ac- and dc-coupled preamplifiers that give a linear response with respect to the infrared light intensity. Data were taken with a temporal resolution of 20 ns, 100 ns, and 10  $\mu\text{s}$  at a spectral resolution of 4  $\text{cm}^{-1}$ . The spectral region covers 3000–800  $\text{cm}^{-1}$  (a combination of the MCT detector cutoff at 800  $\text{cm}^{-1}$  and the bandpass germanium filter cutoff at 3000  $\text{cm}^{-1}$ ) in the time domain of 20 ns to 1 ms. Transient signals were co-added five to six times to improve the signal-to-noise ratio. The pulsed excitation source is the output at 532 nm from a Nd:YAG laser (Quantum-Ray DCR-3, Mountain View, CA), with laser energy of 3 mJ/pulse at a repetition rate of  $\sim 4$  Hz, which is triggered externally through the Nicolet Omnic software. All the measurements were carried out at room temperature.

## RESULTS

Fig. 1 shows the time-resolved FTIR difference spectrum between the photointermediate spectrum formed at a delay time of 600 ns after laser excitation and that of bR ground state, in the 3000–1000- $\text{cm}^{-1}$  frequency range. In the time-resolved difference spectra, the positive bands originate from the formation of photointermediates and the negative bands result from the ground-state depletion as bR mole-

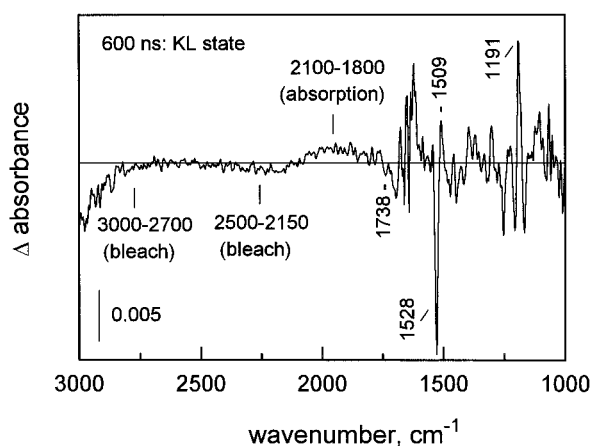


FIGURE 1 Time-resolved FTIR difference spectrum for bR in  $\text{H}_2\text{O}$  in the frequency range of  $3000\text{--}1000\text{ cm}^{-1}$ , measured at room temperature. The spectrum is taken at a delay time of 600 ns, showing the difference between the spectrum of bR and KL state. Continuum band features are observed in the frequency regions of  $3000\text{--}2700\text{ cm}^{-1}$  as a bleach,  $2500\text{--}2150\text{ cm}^{-1}$  as a bleach, and  $2100\text{--}1800\text{ cm}^{-1}$  as a transient absorption.

cules absorb light and go through the photocycle. The spectrum at the 600-ns delay corresponds to the transition stage from bR to the K (or KL) intermediate. The sharp band positions and relative intensities of the bands in this spectrum, such as  $1528$ ,  $1509$ , and  $1191\text{ cm}^{-1}$ , are found to be in good agreement with the  $600\text{--}800\text{-ns}$  averaged spectrum in previous published reports (Sasaki et al., 1993; Wolfgang et al., 1996). The noise in the  $1700\text{--}1540\text{-cm}^{-1}$  region is due to the overlap with the intense water-bending vibration absorption.

Three broad continua are shown in the  $3000\text{--}1800\text{-cm}^{-1}$  frequency region. Two bleach continua are observed, one in the  $2500\text{--}2150\text{-cm}^{-1}$  region and the other one in the  $3000\text{--}2700\text{-cm}^{-1}$  region. One transient absorption continuum band is observed in the  $2100\text{--}1800\text{-cm}^{-1}$  region. In the spectral region of  $2800\text{--}2200\text{ cm}^{-1}$ , a weak and broad continuum band has been observed previously (Olejnik et al., 1992) in the trapped K and L minus bR difference spectra obtained at low temperature. These authors attributed such a broad band to the formation of at least one hydrogen bond in which the proton is not well localized. A recent report (Kandori et al., 1998), however, showed only a broad continuum bleach band in the  $3000\text{--}2700\text{-cm}^{-1}$  region in the K minus bR difference spectrum obtained at 77 K, but not in the  $2800\text{--}2200\text{-cm}^{-1}$  region. Our results show that both absorption and bleach bands in this frequency region can be observed at room temperature by using time-resolved step-scan FTIR. In addition, it should be noted that a continuum absorption band is shown in the  $2100\text{--}1800\text{-cm}^{-1}$  region. In a recent report (Wolfgang et al., 1996), a continuum absorption band feature in the  $1850\text{--}1800\text{-cm}^{-1}$  has also been observed in their  $20\text{--}940\text{-ns}$  spectra. However, no nanosecond or microsecond transient spectra in the

frequency region above  $2000\text{ cm}^{-1}$  have ever been reported so far.

Fig. 2 shows the observed spectra in the  $3000\text{--}1000\text{-cm}^{-1}$  region that are taken at selected delay times (from 50 ns to  $950\text{ }\mu\text{s}$ ) after bR has been photolyzed with 8-ns laser pulses at 532 nm. It should be noted that in both Figs. 1 and 2, the spectral sensitivity has been increased to observe the weak continuum. As a result, the strong bands at  $1528\text{ cm}^{-1}$  and  $1191\text{ cm}^{-1}$  are off scale. It is found that the transient absorption in the  $2100\text{--}1800\text{-cm}^{-1}$  region is observed at 50 ns (Fig. 2 a) as well as  $10\text{ }\mu\text{s}$  (Fig. 2 c). At this time, the positive band at  $1763\text{ cm}^{-1}$  (resulting from the  $\text{C=O}$  stretching mode of the protonated proton acceptor Asp85) appears, whereas the positive band at  $1191\text{ cm}^{-1}$  showed a significant decrease in intensity relative to that in the 50-ns spectrum.

Fig. 3 compares the time dependence of the transient continua at  $2100\text{--}1800\text{ cm}^{-1}$ ,  $2500\text{--}2150\text{ cm}^{-1}$ , and  $3000\text{--}2700\text{ cm}^{-1}$  in the  $20\text{-ns}$  to  $1\text{-}\mu\text{s}$  time domain, with that of

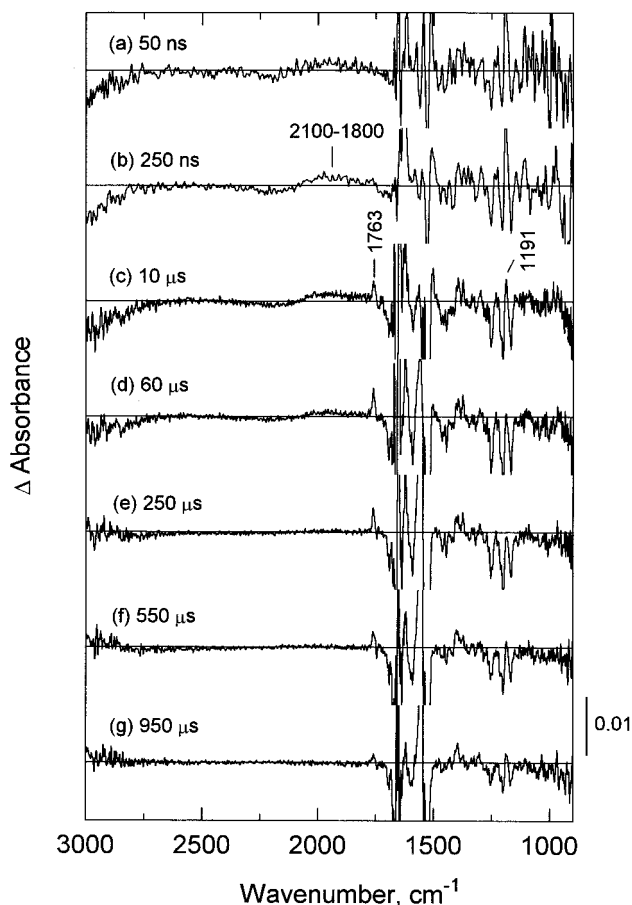


FIGURE 2 Time-resolved FTIR difference spectra for bR in  $\text{H}_2\text{O}$  in the frequency range of  $3000\text{--}1000\text{ cm}^{-1}$ , measured at room temperature, and at various delay times as indicated: (a) 50 ns; (b) 250 ns; (c)  $10\text{ }\mu\text{s}$ ; (d)  $60\text{ }\mu\text{s}$ ; (e)  $250\text{ }\mu\text{s}$ ; (f)  $550\text{ }\mu\text{s}$ ; (g)  $950\text{ }\mu\text{s}$ . This covers the time domain of the KL-L-M portion of the bR photocycle.

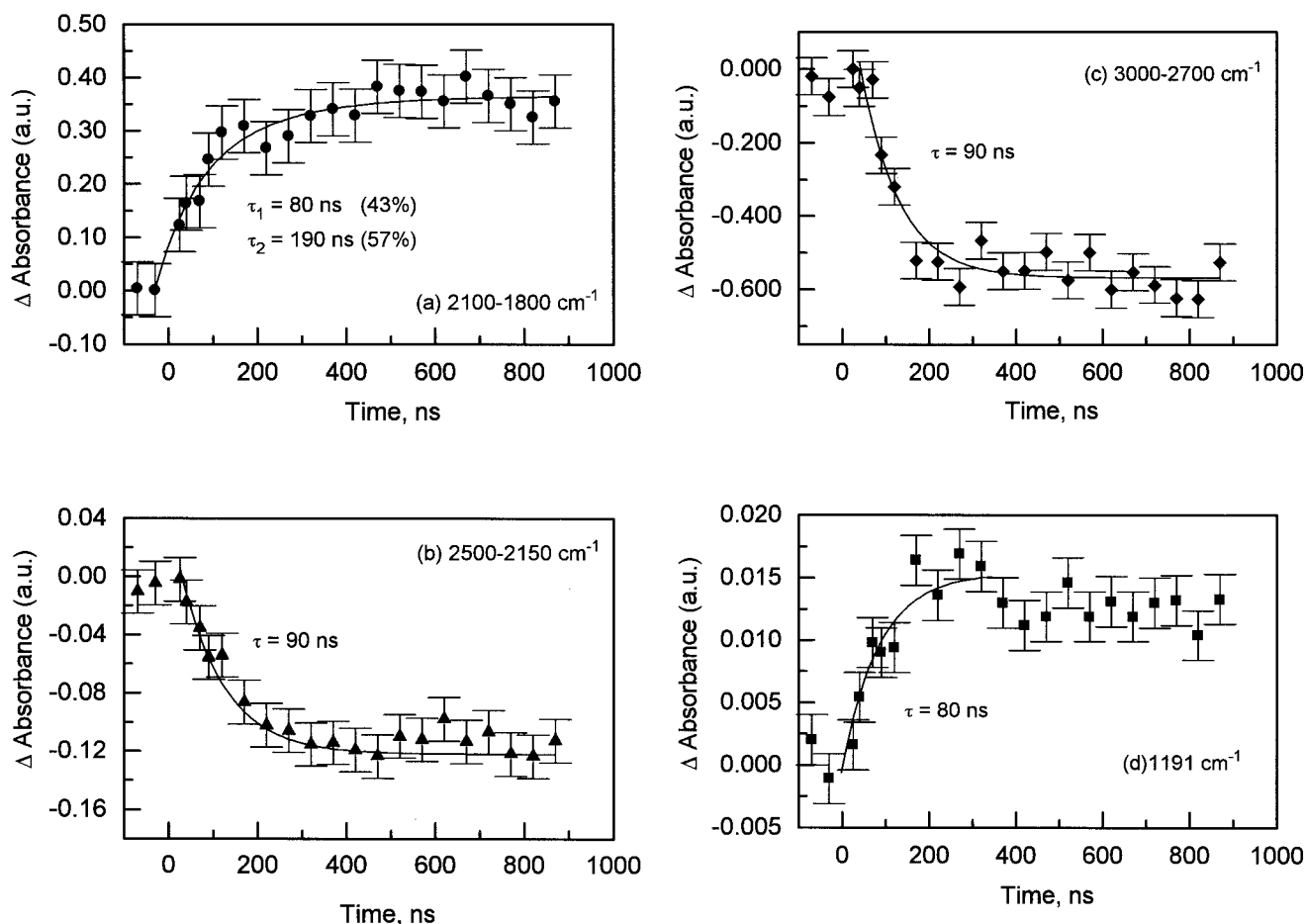


FIGURE 3 The initial temporal behavior of selected bands in the time domain of 20 ns to 1  $\mu$ s. (a) The absorption continuum band in the 2100–1800- $\text{cm}^{-1}$  region; (b) The bleach continuum 2500–2150- $\text{cm}^{-1}$  region; (c) The bleach continuum 3000–2700- $\text{cm}^{-1}$  region; (d) The retinal C—C stretching band at 1191  $\text{cm}^{-1}$ . The fitted results are also shown, with the corresponding rise times given as:  $<80$  ns; 190 ns for formation of the 2100–1800- $\text{cm}^{-1}$  continuum absorption band;  $<90$  ns for formation of the 2500–2150- $\text{cm}^{-1}$  continuum bleach band; and  $<80$  ns for the formation of the 1191- $\text{cm}^{-1}$  band. The 80–90-ns short components observed correspond to the rise time of the detector suggesting that the initial component of these changes corresponds to the formation of the K intermediate.

the sharp positive band at 1191  $\text{cm}^{-1}$  that is due to the 13-*cis* retinal C—C stretching in the K intermediate (Smith et al., 1986). Fitting the results gives two exponential rise components for the absorption continuum band at 2100–1800  $\text{cm}^{-1}$ , with apparent rise times of  $<80$  ns and 190 ns, respectively. The fast component is on the same time scale as that of the formation of the bleach bands ( $<90$  ns; Fig. 3, *b* and *c*) and that of the 1191- $\text{cm}^{-1}$  band ( $<90$  ns; Fig. 3 *d*). It is known that the photoisomerization of retinal occurs in less than half a picosecond (Nuss et al., 1985; Petrich et al., 1987; Mathies et al., 1988; Van den Berg et al., 1990). Because the fingerprint band at 1191  $\text{cm}^{-1}$  is associated with the K formation (Smith et al., 1986), its rise time should be  $\sim 3$  ps instead of  $\sim 90$  ns that we observed here. In fact, the 50–90-ns rise-time component appearing in all these spectra is determined by the rise time of the MCT detector used in the step-scan FTIR spectrometer. Our data thus indicate that the formation of the continuum bands is

possibly associated with the retinal photoisomerization and the K formation. From Figs. 2 and 3, it is shown that the intensity of the observed continua is still developing in the time domain between 100 ns and a few microseconds (within which L is being formed), whereas the intensity of the retinal configuration sensitive band at 1191  $\text{cm}^{-1}$  begins to decay within hundreds of nanoseconds (within which KL begins to decay (Wolfgang et al., 1996)), suggesting that these continua cannot be assigned to retinal vibrations.

Fig. 4 compares the decay of the absorption continuum band in the 2100–1800- $\text{cm}^{-1}$  region (Fig. 4 *a*), the bleach continuum in the 2500–2150- $\text{cm}^{-1}$  region (Fig. 4 *b*), the bleach continuum in the 3000–2700- $\text{cm}^{-1}$  region (Fig. 4 *c*), and the transient absorption of the carbonyl band at 1763  $\text{cm}^{-1}$  (Fig. 4 *d*) (which reflects the protonation of Asp85). The results show that the absorption continuum band in the 2100–1800- $\text{cm}^{-1}$  region has two decay components, with lifetimes of  $\sim 60$  and 300  $\mu$ s. The fast component is similar



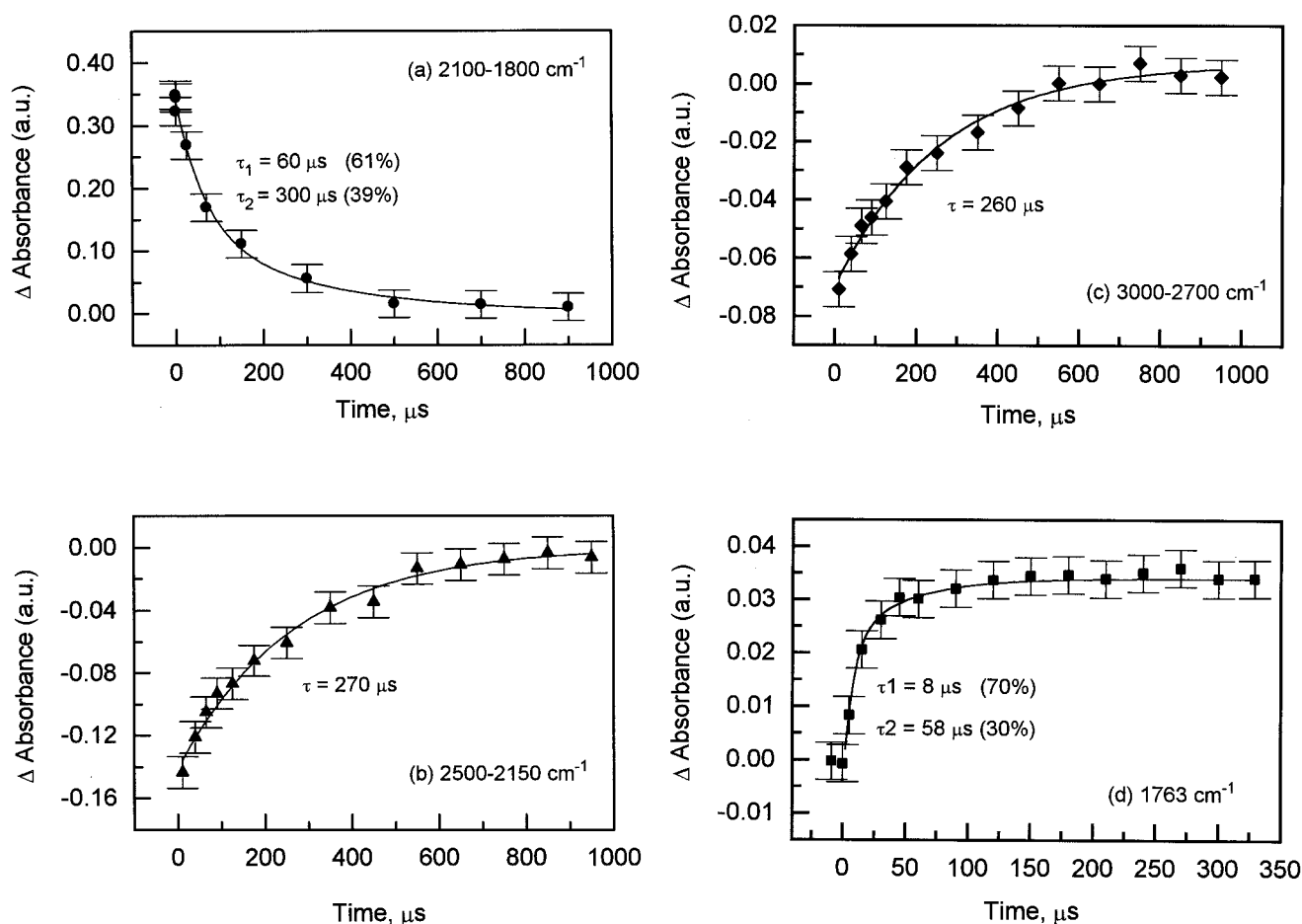


FIGURE 4 The long-time temporal behavior of selected bands in the time domain of micro- to milliseconds. (a) The decay of the transient absorption continuum band in the region of 2100–1800  $\text{cm}^{-1}$ ; (b) The recovery of the bleach continuum band in the region of 2500–2150  $\text{cm}^{-1}$ ; (c) The recovery of the bleach continuum band in the region of 3000–2700  $\text{cm}^{-1}$ ; (d) The rise of the C=O stretching mode absorption of Asp85 at 1763  $\text{cm}^{-1}$ , an indicator of the protonation of Asp85. The fit to the observed results are also shown, with lifetimes given below: 60  $\mu\text{s}$  and 300  $\mu\text{s}$  for the decay of the 2100–1800- $\text{cm}^{-1}$  band; 260–270  $\mu\text{s}$  for the bleach recovery of the BC bands. The formation of the 1763- $\text{cm}^{-1}$  band has two time constants of  $\sim 10 \mu\text{s}$  and  $\sim 60 \mu\text{s}$ , as also reported previously (Wang and El-Sayed, 2000).

to the slow rise-time component of the 1763- $\text{cm}^{-1}$  band, (i.e., of the protonation of the proton acceptor, Asp85), which is found to be  $\sim 60 \mu\text{s}$ , as also seen in previous reports (Siebert et al., 1982; Wang and El-Sayed, 2000). On the other hand, the decay of the two bleach continua has a lifetime of  $\sim 260 \mu\text{s}$ , which is similar to the decay process of the long-lived component of the transient absorption continuum in the 2100–1800- $\text{cm}^{-1}$  region ( $\sim 300 \mu\text{s}$ ), suggesting that these two continua belong to the same polarizable proton chain, or to two polarizable proton systems that are strongly coupled.

Deuterium effects on the observed broad continua band intensity have also been investigated in the present study. Fig. 5 shows a series of transient spectra of bR in  $\text{D}_2\text{O}$  solvent, at delay times of 200 ns, 600 ns, and 10  $\mu\text{s}$ , respectively. It is obvious that there is no significant continuum absorption or bleach in  $\text{D}_2\text{O}$  observed in these regions as observed in  $\text{H}_2\text{O}$  (above 2700  $\text{cm}^{-1}$  and in the

regions of 2500–2150  $\text{cm}^{-1}$  and 2100–1800  $\text{cm}^{-1}$ ). The distorted baseline and spectral noise in the 2750–2000- $\text{cm}^{-1}$  spectral region are due to the strong absorption of the D—O stretching vibration mode of  $\text{D}_2\text{O}$ . This absorption could obscure the much weaker continua due to the polarizable deuterons. These results indicate that the observed continua in the bR/ $\text{H}_2\text{O}$  system show a deuterium isotopic effect and, therefore, involve proton oscillation. The deuterium isotopic effect has been previously reported in the 2100–1800- $\text{cm}^{-1}$  region (Wang and El-Sayed, 2000).

## DISCUSSION

### The polarizable hydrogen-bonded continua

In a series of experiments and theoretical calculations, Zundel and co-workers have demonstrated that delocalized protons in hydrogen-bonded systems have several broad con-

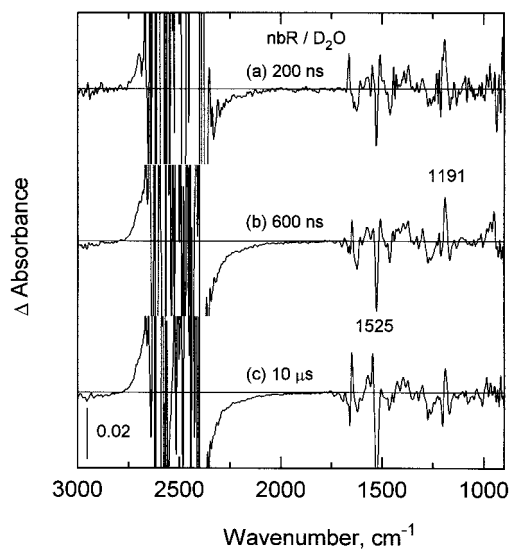


FIGURE 5 Time-resolved FTIR difference spectra in the bR photocycle in  $D_2O$  in the frequency range of  $3000\text{--}1000\text{ cm}^{-1}$ , measured at room temperature, at delay times as follows: (a) 200 ns; (b) 600 ns; (c) 10  $\mu\text{s}$ . The spectra show that the continua observed in  $H_2O$  at  $>2700\text{ cm}^{-1}$ ,  $2500\text{--}2150\text{ cm}^{-1}$ , and  $2100\text{--}1800\text{ cm}^{-1}$  (Fig. 2) are not observed in  $D_2O$ . This strongly suggests that the continuum absorption in the  $3000\text{--}1800\text{ cm}^{-1}$  band results in from the oscillating motion of protons (The strong noise-like signal between  $2750$  and  $2200\text{ cm}^{-1}$  is due to the absorption of  $D_2O$  in this region).

tinuum bands in the mid-infrared region (Zundel and Brzezinski, 1992). In such a hydrogen-bonded network, a large proton polarizability results from the delocalization of proton into a hydrogen-bonded chain, which gives rise to an absorption continuum. These researchers showed that such continua are observed when the hydrogen-bonded dimer is homogeneous ( $H\cdots A\cdots H\cdots A$ ) or heterogeneous ( $H\cdots A\cdots H\cdots B$ ). It is also observed between hydrogen bonds involving cations (e.g.,  $\text{—N}^+\cdots H\cdots C\text{—}R_2$ ) and anions (e.g.,  $\text{=C—O}^-\cdots H\cdots O\text{—}R$ ).

It should be mentioned that a polarizable hydrogen-bonded proton system has continua in more than one region with different intensities. Thus, during the photocycle of bR, as the polarizable hydrogen-bonded systems are changing, the intensity of the continua spectra for different intermediates changes in different regions by different amounts. Therefore, what we observe in each region at different times is the net change of the bleach due to the disappearance of the photolyzed bR and of the absorption due to the appearance of intermediates on the time scale examined. The observation of a net transient continuum absorption in a certain region of the spectrum signals the formation of one or more polarizable hydrogen-bonded systems that are more intense than those of the bleached one(s) in bR ground state.

The present results show a net absorption continuum in the  $2100\text{--}1800\text{ cm}^{-1}$  region with two rise times of  $<80$  and  $190\text{ ns}$  and two decay components of  $60$  and  $300\text{ }\mu\text{s}$  (see

Figs. 3 and 4). Previous studies (Riesle et al., 1996; Rammelsberg et al., 1998; Wang and El-Sayed, 2000) showed a transient bleach (instead of absorption) in the frequency region of  $1900\text{--}1800\text{ cm}^{-1}$  that is formed within hundreds of microseconds. In the present study, the decay time determined for the transient absorption in the  $2100\text{--}1800\text{ cm}^{-1}$  region is in agreement with the previously reported rise times for the formation of the observed bleach band in the  $1900\text{--}1800\text{ cm}^{-1}$  region ( $50\text{--}70\text{ }\mu\text{s}$ ). However, there is an apparent difference in the assignment of spectral characteristics (absorption versus bleach). We believe the reason for this is the baseline correction that is related directly to the infrared bandpass filter used in the step-scan mode of FTIR. In the present study, we use a bandpass filter that covers the entire  $3000\text{--}500\text{ cm}^{-1}$  range, which is much broader than most of the band filters used in the previous studies (Rammelsberg et al., 1998; Wang and El-Sayed, 2000) that cover only the  $1900\text{--}900\text{ cm}^{-1}$  region. Therefore, a decay process of a small positive transient absorption signal at a certain frequency observed in the broad spectral region could be seen as a formation process of a small negative transient bleach signal at the cutoff frequency portion of a narrow bandpass filter.

### The molecular origin of the polarizable hydrogen-bonded proton continua in bR

The net bleach continuum spectra in the  $2500\text{--}2150\text{ cm}^{-1}$  region and above  $2700\text{ cm}^{-1}$  appear upon the formation of K and decay in  $300\text{ }\mu\text{s}$ . Thus, one or more polarizable hydrogen-bonded protons present in bR change upon retinal photoisomerization. This polarizable hydrogen-bonded system must be within the active site and very close to the retinal molecule. Thus, beside the observed bleach continuum, the formation of K is accompanied by the formation of one or more polarizable hydrogen-bonded systems, which gives the observed net absorption in the  $2100\text{--}1800\text{ cm}^{-1}$  region. This continuum has two rise times of  $<80$  and  $190\text{ ns}$  and two decay times of  $50\text{--}60\text{ }\mu\text{s}$  and  $300\text{ }\mu\text{s}$ . Because the component that rises in  $<80\text{ ns}$  and decays in  $300\text{ }\mu\text{s}$  has similar kinetics as that of the observed bleach continua, one is tempted to assign it to the polarizable hydrogen-bonded system resulting from the photoisomerization. The other component that rises in  $190\text{ ns}$  and decays in  $50\text{--}60\text{ }\mu\text{s}$  could thus be another polarizable hydrogen-bonded system that is involved in pumping the proton to the surface. This is because the L-to-M transition occurs in  $\sim 50\text{--}60\text{ }\mu\text{s}$ . It is also on the same time scale as the appearance of the pumped proton to the surface (Heberle and Dencher, 1992; Alexiev et al., 1995) and that for the deprotonation of Tyr185, as determined optically (Dupuis et al., 1985) and from FTIR difference spectroscopy (Rothschild et al., 1990a).

From the x-ray studies (Edman et al., 1999; Luecke et al., 1999b), one can propose two possible polarizable hydrogen-bonded systems in bR. The first one is between the PSB and

a water molecule (water 402) ( $\text{N}^+ - \text{H} \cdots \text{OH}_2$ ). The second one is between Tyr185 and Asp212 ( $-\text{Ph}-\text{O}-\text{H} \cdots \text{O}^-_2\text{C}-$ ). They could form either two independent hydrogen-bonded systems or just one system. For example, the one system could involve the PSB, water 402, Asp85, Tyr185, Asp212, and water 401, as concluded from the recent x-ray structure of bR (Luecke et al., 1999a). In K, there is no doubt that this stable hydrogen-bonded structure is disturbed upon retinal photoisomerization. There are several independent studies whose results suggest a large change in the charge distribution takes place in the bR-K transition. Movement of a water molecule occurs (Edman et al., 1999) and Tyr185 seems to become partially deprotonated (Rothschild et al., 1990a), and a rapid charge separation is detected (Keszthelyi and Ormos, 1983; Trissl, 1990; Liu and Ebrey, 1988; Liu, 1990) that has a rise time of  $<30$  ps (Groma et al., 1984). Thus, the changes in the nature of the polarizable hydrogen bonds upon K formation might be responsible for the appearance of the net absorption in the  $2100\text{--}1800\text{-cm}^{-1}$  region and the observed net bleach in the region of  $2500\text{--}2150\text{ cm}^{-1}$  and above  $2700\text{ cm}^{-1}$ .

The FTIR difference spectrum between bR and trapped K (Rothschild et al., 1990a) showed that the degree of Tyr185 protonation changes, as monitored by the  $1277\text{-cm}^{-1}$  and  $1272\text{-cm}^{-1}$  bands that are absent in the Asp212 mutant. This led them to conclude that a polarizable hydrogen bond exists between Tyr185 and Asp212 in bR, which leads to the shift of Tyr185 hydrogen toward Asp212 upon K formation. One might then conclude that this contributes to the net bleach and net absorption continua during the bR-to-K transformation. If indeed the two polarizable hydrogen-bonded systems (the  $\text{PSB} \cdots \text{H}_2\text{O}(402)$  and  $\text{Tyr185} \cdots \text{Asp212}$ ) change during the formation of K, one can assign the fast component in the absorption continuum ( $<90$  ns) to changes in the  $\text{PSB} \cdots \text{H}_2\text{O}(402)$  system and the slow component (190 ns) to changes in the  $\text{Tyr185} \cdots \text{Asp212}$  system. This is because one might expect that the structure around the  $\text{C}=\text{N}^+\text{H}$  moiety of the PSB changes on the same time scale as the photoisomerization time scale (0.5 ps).

Optical spectroscopy has also shown that during the M formation, some tyrosine residues deprotonate on the same time scale as the deprotonation of PSB (Scherrer et al., 1981; Lemke and Oesterhelt, 1981; Lam et al., 1983; Dupuis and El-Sayed, 1985). This tyrosine residue has been shown later to be Tyr185 (Dunach et al., 1990). The fact that the absorption continuum in the  $2100\text{--}1800\text{-cm}^{-1}$  region appears with a slow component of 190 ns and its short component disappears upon M formation strongly suggests our assignment that the polarizable hydrogen bond between Asp212 and Tyr185 makes a good contribution to it. Upon K formation a displacement of Tyr185 hydrogen of the hydroxyl group occurs in the  $\text{Tyr185} \cdots \text{Asp212}$  hydrogen-bonded system. In the L-to-M transition, this hydrogen is ionized completely. The recovery of the ground-state bleach in  $300\text{ }\mu\text{s}$  cannot be due to a polarizable hydrogen-bonded

proton involving the PSB or the tyrosinate because both of these recover on a much longer time scale (outside our limit). It must then involve water molecules that do not affect the absorption of the retinal or the tyrosinate ion. Of course, it could also result from the formation of a new polarizable hydrogen-bonded proton system whose intensity compensates that of the bleach of the ground state in this frequency region.

### The polarizable hydrogen-bonded chains and the proton pump mechanism

It has been found that the Schiff base proton is transferred to Asp85 but remains bound to it until the recovery of bR at the end of the photocycle (Siebert et al., 1982; Bousche et al., 1992; Souvignier and Gerwert, 1992). Another proton appears on the surface on the same time scale as the protonation of Asp85 (Heberle and Dencher, 1990, 1992; Scherrer et al., 1992; Heberle et al., 1993). Active research has been carried out in the past decade aimed at identifying the molecular origin of the proton being pumped to the surface and thus the mechanism of the proton pump process itself. The proton-releasing process has been proposed to involve Arg82 (Otto et al., 1990; Balashov et al., 1993, 1995; Govindjee et al., 1996), Glu204 (Misra et al., 1997; Kandori et al., 1997; Richter et al., 1996; Brown et al., 1995, 1996), Glu194 (Dioumaev et al., 1998), and Arg134 (Lu et al., 2000) residues. Earlier, other proposals suggested that an ionized internal water molecule could be the source of the surface protons (El-Sayed et al., 1995).

The transport of the proton to the surface via a chain of water molecules was proposed as early as 1981 (Nagle and Mille, 1981), long before any spectroscopic experimental evidence of the presence of internal water molecules was reported. The presence of internal water molecules in the active site of bR has been detected by using Raman spectroscopy (Hildebrandt and Stockburger, 1984; Deng et al., 1994; Althaus et al., 1995), by FTIR (Chang et al., 1991; Maeda et al., 1992; Fischer et al., 1994), by neutron diffraction (Papadopoulos et al., 1990), by nuclear magnetic resonance (NMR) (Lugtenburg et al., 1988), and recently by x-ray diffraction (Pebay-Peyroula et al., 1997; Luecke et al., 1999b).

Recently, a proton pump mechanism was proposed (Luecke et al., 1999a) from the observed changes in the x-ray diffraction upon M formation. In this model, the changes in the  $\text{pK}_a$  of different residues resulting from changes in the water structure within the hydrophobic interior of the protein is proposed to lead to the final deprotonation of Glu204. There are, however, a number of facts for which this model does not account. 1) Time-resolved FTIR studies showed that Glu204 is not deprotonated in M. 2) Tyr185 is deprotonated in M but not in this model. 3) Both the PSB and Tyr185 deprotonate at the same time as the appearance of the proton on the surface. And 4) three of the seven internal

water molecules disappear upon the M formation (water 402, 406, and 403).

Because Glu204 is not deprotonated in M, there must be another source for the proton appearing on the surface in M. From the above observations, one concludes that the source of the proton that appears on the surface could either come from Tyr185 or from one of the disappearing water molecules. Tyr185 can give up its proton to initiate a rapid transport proton on the hydrogen-bonded chain made of water and other residues to reach the surface. It needs a trigger for its complete deprotonation process that should be turned on in about the same time as that for the proton transfer from the PSB to Asp85 via water 402. If as discussed above, there is a strongly coupled hydrogen-bonded network involving the two hydrogen-bonded polarizable protons (that of the PSB and that of Tyr185), then the trigger could simply be the motion of one or more water molecules. This could decrease the pKa of both the PSB and Tyr185 and/or decrease the pKb of the conjugated base of Asp85 and Asp212 acids. Theoretical calculation is needed to examine this possible mechanism.

It was previously suggested (El-Sayed, 1992), that the trigger could involve an interior cation, stabilized by a cluster of two to three water molecules and  $O^-$  of the carboxylate groups of the amino acid residues. In the L-M transition, a simultaneous exposure of the PSB and Tyr185 to the positive charge on this cation can take place (Dupuis et al., 1985). This greatly reduces the pKa values of both the PSB and Tyr185 as a result of the repulsion between the cation and the PSB and the attraction between the tyrosinate anion and the cation. This model offers a switch that ionizes the PSB proton and the Tyr185 proton on a comparable time scale. The PSB proton is transferred to Asp85 and the Tyr185 proton is transferred to the surface via a hydrogen-bonded water chain.

It was also proposed previously (El-Sayed et al., 1995) that Asp85 might be one of the ligands that stabilizes the metal cation (together with water 402 and other water molecules). As it gets neutralized, one of the water molecules (e.g., water 402) ionizes, giving rise to one proton and one  $OH^-$ . The  $OH^-$  group substitutes for Asp85 as a ligand to stabilize the metal cation. The resulting proton could thus be transferred to the surface via the hydrogen-bonded chains present in the system. In this model, the metal cation switch might even release two protons, one from Tyr185 and another one from a water molecule, only one of them will be transferred to the surface via the hydrogen-bonded network shown in the x-ray studies (Edman et al., 1999; Luecke et al., 1999a,b).

As attractive as it is, the above model requires a metal cation to be bound within the reaction center that is ligated to water molecules, and few amino acid residues. This metal cation complex can communicate effectively with the surface through the hydrogen-bonded network. There are many references in the literature supporting the presence of one or

two specific interior cation-binding sites (Chang et al., 1985; Dupuis et al., 1985; Chronister et al., 1986; Arikai and Lanyi, 1986; Corcoran et al., 1987; Dunach et al., 1987; Jonas et al., 1990; Mitra and Stroud, 1990; Zhang et al., 1992, 1993; Yang and El-Sayed, 1995; Stuart et al., 1995; Yoo et al., 1995; El-Sayed et al., 1995; Tan et al., 1996; Birge et al., 1996) as well as those supporting the presence of only surface binding sites for all metal cations (Szundi and Stoerkenius, 1987, 1988, 1989; Fu et al., 1997; Tuzi et al., 1999; Varo et al., 1999). The fact that the recent x-ray structures of bR confirm the presence of a number of internal bound water molecules (Edman et al., 1999; Luecke et al., 1999a) makes the proposal of having an interior metal cation not out of the question. Furthermore, the presence of a hydrogen-bonded network connecting the extracellular surface to the reaction center is expected to keep the reaction center with its PSB and Asp85 in rapid equilibrium with the surface potential. This might eliminate the most serious results against the internal binding sites that are based on the observations correlating the surface potential to the protonation state of Asp85 and the color change of the retinal. Previous theoretical calculations (Tan et al., 1996; Pardo et al., 1998) have already shown the possibility of having a metal cation within the reaction center. With the recent three-dimensional atomic structure data of bR, now known exactly, a more accurate *ab initio* calculation can be used to find out whether or not a cation can bind within the reaction center of bR.

## Possible conclusions

Combining previous spectroscopic and x-ray studies and the present results one reaches the following possible conclusions. 1) More than one hydrogen-bonded chain exists in bR that changes upon retinal photoisomerization. 2) One of them involves the PSB- $H_2O$ (402)-Asp85 and is responsible for the transfer of the proton from the PSB to Asp85 through the water molecule during M formation. 3) Another chain involves the Tyr185 and Asp212. 4) These two chains could also be one large chain involving PSB- $H_2O$ (402), Asp85, Tyr185, Asp212, and  $H_2O$ (401) and other water molecules (e.g., 406). 5) The previously observed similar deprotonation rate of the PSB and Tyr185 could be explained by the presence of one trigger. 6) This trigger could be the photoisomerization process followed by  $H_2O$  rearrangements of the water molecules around the PSB and Tyr185 (and Asp212) that leads to changes in their pKa values. 7) This trigger could also be the approach and interaction of an interior metal cation with both the PSB and Tyr185. 8) The source of the pumped surface proton could thus be Tyr185 and/or an ionizable  $H_2O$  molecule ligated to the metal cation.

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